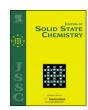
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# Preparation and structural study from neutron diffraction data of Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub>

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#### ABSTRACT

The title compound has been prepared as polycrystalline powder by thermal treatments of mixtures of  $Pr_6O_{11}$  and  $Pr_6O_{12}$  and  $Pr_6O_{13}$  and  $Pr_6O_{13}$  are reported here for the first interesting catalytic properties, but its true stoichiometry and crystal structure are reported here for the first time. It is cubic, isostructural with  $Pr_6O_{13}$  and  $Pr_6O_{13}$ 

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## 1. Introduction

Metal oxides simultaneously containing transition metals and lanthanide ions present a panoply of interesting properties and applications, spanning from superconductivity (e.g. in  $(La,Sr)CuO_{4-\delta}$ ) [1] to colossal magnetoresistance (e.g. in  $(La,Ca)MnO_3$ ) [2] or catalysis (e.g.  $CaFeO_{3-\delta}$ ) [3]. When both metal ions posses a variable oxidation state, the prospect of charge transfer between them appears as an attractive possibility, perhaps associated with novel phenomena and functionalities. As a requirement, the chemical potential of the f ion should be energetically close to the Fermi level of the transition-metal d-bands [4]. This circumstance is accomplished when the f ion is a light lanthanide, particularly Ce or Complex or Complex where the Complex or Complex

Within the field of heterogeneous catalysis, particularly appealing are the ternary Pr–Mo–O [7] and quaternary Bi–Pr–Mo–O [8–10] phases, which are active as catalysts for the partial oxidation of propene and isobutene. An oxide named  $Pr_2MoO_6$  is present in the phase diagram of the Pr–Mo–O system and participates in the catalytic activity for selective olefin oxidation [11–13]. This  $Pr_2MoO_6$  oxide has also been reported among the other members of the well-described  $R_2MoO_6$  family (R=rare earths) [14]. Depending on the size or the synthetic conditions,  $R_2MoO_6$  compounds have been described to crystallize in three polymorphs, with monoclinic ( $\alpha$ ), cubic ( $\beta$ ) and tetragonal ( $\gamma$ ) symmetries [14]. Most of the small-sized rare-earth cations form the monoclinic ( $\alpha$ ) phase, which is closely

related to the scheelite-type structure. Scheelite itself is a super-structure of fluorite (CaF<sub>2</sub>, cubic, a=5.5544 Å, Fd-3m) [15]. The crystal structure of the  $\alpha$ -phase has been recently refined for Tb<sub>2</sub>MoO<sub>6</sub> in the I2/c space group from neutron powder diffraction data [14], as well as those of the remaining  $\alpha$ - $R_2$ MoO<sub>6</sub> polymorphs for small-sized rare earths, R=Y, Dy, Ho, Er, Tm, Yb [16]. The La polymorph is tetragonal (space group  $I4_1/acd$ ,  $\gamma$  phase) [14] whereas the Ce, Pr and Nd compounds have been described as pseudocubic ( $\beta$  phase) [4,14] although the space group symmetry and structural features have not been reported so far. For R=Pr, Nd it seems that the materials prefer either the pseudocubic or the tetragonal structures, depending on the synthesis conditions: the tetragonal form is the high-temperature phase, whereas the pseudocubic structure ( $\beta$ ) can be stabilized at low temperatures [14].

In this work we have focused on the actual composition and crystal structure of the so-called  $Pr_2MoO_6$  pseudocubic oxide. In fact it presents a stoichiometry  $Pr_5Mo_3O_{16}$  and it consists of a cubic superstructure of fluorite, where a mixed or intermediate valence for Pr and/or Pr Mo is expected from the nominal composition. Given the weak Pr-ray scattering power of oxygen compared to that of praseodymium or molybdenum atoms, neutron powder diffraction has been the technique of choice for this investigation. A bond-valence study was of much help to unveil the valence state for the involved metals in this oxide.

#### 2. Experimental

A sample of nominal composition  $Pr_2MoO_6$  was prepared by a solid state reaction from stoichiometric amounts of analytical

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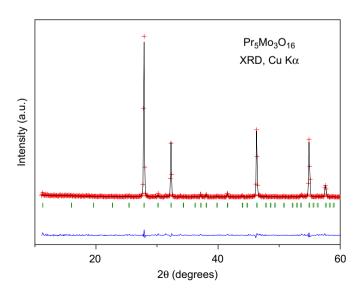
E-mail address: ja.alonso@icmm.csic.es (J.A. Alonso).

grade  $Pr_6O_{11}$  and  $MoO_2.$  The reactants were mixed, ground and calcined at 1100  $^{\circ}\text{C}$  for 12 h in air.

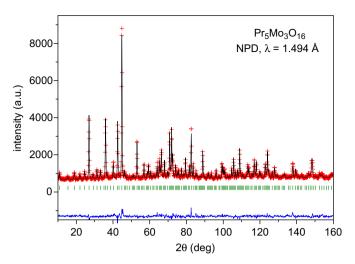
The initial characterization of the product was carried out by laboratory X-ray diffraction (XRD) (Cu  $K\alpha$ ,  $\lambda = 1.5406$  Å). A roomtemperature neutron powder diffraction (NPD) diagram was collected at the HRPT diffractometer of the SINQ spallation neutron source at PSI, with a wavelength of 1.494 Å. The high intensity mode was used; the collection time was 2 h. The structure has been solved ab-initio with the FOX [17] software. The refinement of the crystal structure was performed by the Rietveld method, using the FULLPROF refinement program [18], with the use of its internal tables for the coherent scattering lengths of corresponding elements. The peak profiles were fitted by the Thompson-Cox-Hastings pseudo-Voigt function corrected for axial divergence asymmetry. The following parameters were refined in the final run: scale factor, background coefficients, zero-point error, pseudo-Voigt corrected for asymmetry parameters, positional coordinates and anisotropic thermal factors.

#### 3. Results and discussion

A polycrystalline oxide of nominal composition Pr<sub>2</sub>MoO<sub>6</sub> was obtained as a pale yellow powder. Fig. 1 shows the corresponding XRD diagram, fitted to the structural model described below. The pattern seems to be characteristic of a fluorite-like structure (cubic. a=5.54 Å). However, the NPD pattern shows many more reflections that cannot be indexed in a simple fluorite subcell, already indicating that Pr and Mo are not randomly distributed over the same crystallographic positions, but are long-range ordered in a fluorite superstructure. On the other hand, this XRD diagram is very similar to that given by Antonio et al. [4] for Ce<sub>2</sub>MoO<sub>6</sub>, for which the crystal structure was not reported. The superstructure reflections were more clearly visible in the NPD diagram, which was indexed with the TREOR programs in a cubic unit-cell with a=11.084 Å. A pattern matching (Le Bail fit) in this unit cell leads to an excellent agreement between observed and calculated profiles. The systematic extinctions were consistent with the Pn-3n space group (No. 222). In the asymmetric unit there are two independent Pr atoms at 12e(x, 1/4, 1/4) and 8c(0,0,0) positions, Mo at 12d(0, 3/4, 1/4) sites and the two kinds of oxygen atoms at 48i(x, y, z)and 16f(x, x, x) positions. Fig. 2 illustrates the goodness of the NPD fit. The crystallographic formula indeed corresponds to Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub>,



**Fig. 1.** XRD pattern of  $Pr_5Mo_3O_{16}$ . The Rietveld fit corresponds to the structural model described in the Pn-3n space group, with a=11.08979(9) Å.



**Fig. 2.** Observed (crosses), calculated (solid line) and difference (at the bottom) NPD profiles for  $Pr_5Mo_3O_{16}$  at 295 K. The tick marks indicate the positions of the allowed Bragg reflections.

**Table 1** Structural parameters for  $Pr_5Mo_3O_{15}$  refined in the cubic Pn-3n space group (No. 222) at room temperature from NPD. Lattice parameters: a=11.08979(9) Å and V=1363.86(2) ų. Discrepancy factors:  $R_p$ =3.46%,  $R_{wp}$ =4.49%,  $R_{exp}$ =3.19%,  $\chi^2$ =1.98 and  $R_{Bragg}$ =7.09%.

Atom	Site	х	у		z	$B_{\rm eq}(\mathring{A}^2)$
Pr1	12e	0.0089(5)	0.2500		0.2500	0.51(8)
Pr2	8 <i>c</i>	0.0000	0.0000		0.0000	1.32(16)
Mo	12d	0.0000	0.7500		0.2500	0.98(7)
01	48i	0.5851(2)	0.3629	9(2)	0.8275(2)	1.74(4)
02	16 <i>f</i>	0.1203(3)	0.1203	3(3)	0.1203(3)	0.71(4)
Atom	$\beta_{11}$	opic thermal f $\beta_{22}$	Factors ( $\times$ 1 $\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
Pr1	4(4)	13(3)	13(3)	0	0	0
Pr2	27(3)	27(3)	27(3)	-4(3)	-4(3)	-4(3)
Mo	29(4)	16(2)	16(2)	0 ` ´	0 `´	0 ` ´
		, ,	42(2)	-9(2)	17(2)	-1(2)
01	41(2)	23(2)				
O1 O2	41(2) 14(1)	23(2) 14(1)	14(1)	0	0	0

Z=4. In fact, this structural arrangement resulted to be isotypic with CdTm<sub>4</sub>Mo<sub>3</sub>O<sub>16</sub> [19]. The refined atomic coordinates after the refinement are listed in Table 1, and the bond distances and angles are given in Table 2.

A representation of the crystal structure is shown in Fig. 3. The oxygen coordination polyhedra for Pr1, Pr2 and Mo are shown in Fig. 4. Both Pr atoms are coordinated to eight oxygen atoms, within rather distorted scalenohedra. For Pr10<sub>8</sub> the polyhedron contains 4 Pr1-O1 bond lengths of 2.589(4) Å, and 4 Pr1-O2 of 2.380(4) Å with average  $\langle Pr-O \rangle$  distances of 2.485 Å; in Pr2O<sub>8</sub> there are 6 Pr2-O1 bond lengths of 2.619(2) Å and 2 Pr2-O2 distances of 2.311(3) Å, with a  $\langle Pr2-0 \rangle$  average of 2.542 Å. We can define the relative distortion of the polyhedra with the  $\Delta_d$  parameter, concerning the deviation of Pr-O distances with respect to the average  $\langle Pr-O \rangle$ value, as  $\Delta_d = (1/8) \sum_{n=1,8} [(d_n - \langle d \rangle)/\langle d \rangle]$  [2]; we obtain  $\Delta_d(\Pr 1) =$  $17.7 \times 10^{-4}$  and  $\Delta_d(Pr2) = 27.7 \times 10^{-4}$ . Mo atoms are bonded to four oxygen atoms at distances of 1.788(2) A conforming quasi-regular tetrahedral units (O-Mo-O angles are 116.3(2)° and 106.2(2)°). A description of the structure can be obtained from Fig. 5. Pr10<sub>8</sub> scalenohedra form trimer units by sharing three edges via O2 atoms; every two of these units form a quasi-cubic cavity where Pr2 atoms are located; thus the Pr2O<sub>8</sub> scalenohedron shares three edges

Table 2 Main bond distances (Å) and selected angles (deg.) for  $Pr_5Mo_3O_{15}$  determined from NPD data at RT.

Pr1-O1 ( × 4) Pr1-O2 ( × 4)	2.589(4) 2.380(4)	
Pr2-O1 ( × 6) Pr2-O2 ( × 2)	2.619(2) 2.311(3)	
Mo-O1 ( × 4) O1-Mo-O1 ( × 2)	1.788(2) 116.3(2)	
01-Mo-01 ( × 2)	106.2(2)	

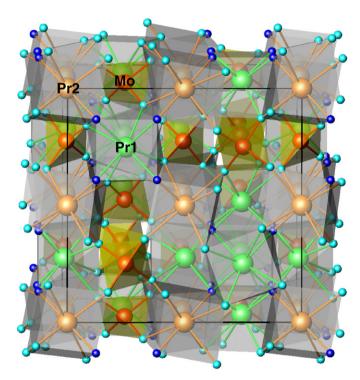


Fig. 3. View of the crystal structure of  $Pr_5Mo_3O_{16}$  along an axis:  $PrO_8$  scalenohedra are long-range ordered with  $MoO_4$  tetragonal units.

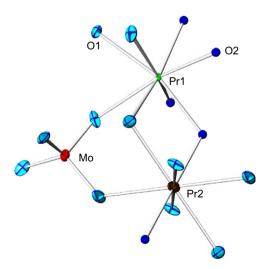
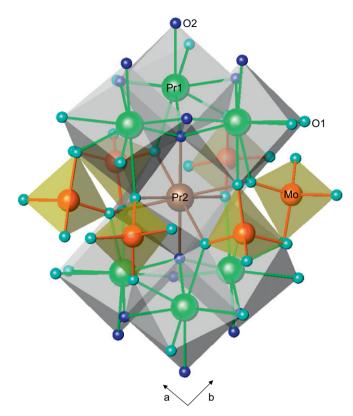


Fig. 4. ORTEP plot of the asymmetric unit of  $Pr_5Mo_3O_{16}$ , showing the oxygen coordination of  $Pr_1$ ,  $Pr_2$  and Mo.



**Fig. 5.** Partial view of the crystal structure of  $Pr_5Mo_3O_{16}$ ;  $Pr1O_8$  scalenohedra share three edges forming a quasi-cubic cavity where Pr2 atoms are located; the  $MoO_4$  tetrahedra are corner connected to both Pr scalenohedra.

with the upper  $Pr_3O_{19}$  unit and three edges with the lower unit. The  $MoO_4$  tetrahedra are bridging the upper and lower units, by sharing corners via the O1 oxygen. The  $MoO_4$  tetrahedra are isolated from each other (they do not share common oxygen), forming discrete units.

It is simple to understand the relationship of this complex superstructure with the basic fluorite CaF<sub>2</sub> structure, containing regular CaF<sub>8</sub> units in a cubic coordination. The unit-cell parameters of the cubic structure of Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub> are related with the simple cubic  $a_f$  fluorite unit cell as  $a \approx 2a_f$  ( $a_f = 5.5 \text{ Å}$ ).  $Pr_5Mo_3O_{16}$  can be rewritten as M<sub>8</sub>O<sub>16</sub>, containing 8 fluorite MO<sub>2</sub> units per formula and 32 MO<sub>2</sub> units per unit cell. The complex cubic superstructure thus results from the long-range ordering between PrO<sub>8</sub> and MoO<sub>4</sub> polyhedra across the crystal. In the fluorite structure every CaF<sub>8</sub> cubic polyhedron shares edges with 8 neighboring units. In Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub> this long-range arrangement is perturbed by the presence of MoO<sub>4</sub> tetrahedra, corner linked with 4 neighboring PrO<sub>8</sub> units. It is important to highlight that this structural type does not contain oxygen vacancies (as it happens in many fluorite-like oxygen-defective oxides, such as Pr<sub>6</sub>O<sub>11</sub>, deriving from PrO<sub>2</sub> by removal of 1/12 of the anions [20]). In the present case, the strong preference of Mo for the tetrahedral coordination gives rise to a deformation of the oxygen sublattice, inducing the presence of cavities adjacent to the MoO<sub>4</sub> tetrahedra, as it is highlighted in Fig. 6. These cavities could provide an easy diffusion path for oxygen ions, via interstitial positions, which is certainly worth exploring; the oxygen transport properties of this material have not been reported up to date.

It is interesting to get an insight into the valence states of the different cations present in the solids by means of the Brown's bond valence model [21,22]. It gives a phenomenological relationship between the formal valence of a bond and the corresponding bond length. Table 3 includes the valences calculated for Pr, Mo and O

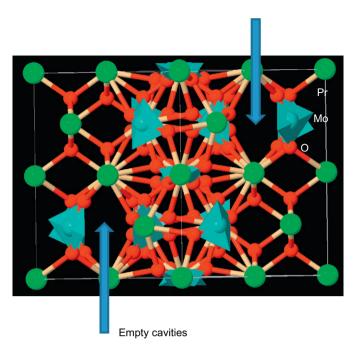


Fig. 6. [1 1 0] projection of the crystal structure, highlighting the presence of cavities adjacent to the  $MoO_4$  tetrahedra.

**Table 3** Valences<sup>a</sup> determined from the individual Pr-O and Mo-O distances in Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub>.

Pr1	3.26(1)
Pr2	2.887(9)
Mo	5.52(2)
01	1.947(9)
02	2.19(1)
GII	0.277
GII	U.Z//

<sup>&</sup>lt;sup>a</sup> The valence is the sum of the individual bond valences  $(s_i)$  for Pr–O and Mo–O bonds. Bond valences are calculated as  $s_i = \exp[(r_0 - r_i)/B]$ ; B = 0.37,  $r_0 = 1.743$  for the Mo<sup>6+</sup>–O<sup>2-</sup> pair and  $r_0 = 2.135$  for the Pr<sup>3+</sup>–O<sup>2-</sup> pair, from Ref. [22]. Individual Pr–O and Mo–O distances  $(r_i)$  are taken from Table 2. The global instability index (GII) is calculated as the root mean of the valence deviations for the j = 1, ..., N atoms in the asymmetric unit, according to  $GII = (\Sigma_j [\Sigma_i(s_{ij} - V_j)^2]/N)^{1/2}$ .

from the individual Pr-O and Mo-O distances in Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub>. The valence of Pr1 ions is significantly higher than the expected value of +3 whereas the valence of Pr2 atom is slightly lower than +3. This result suggests that Pr1 atoms are overbonded while Pr2 are slightly underbonded in this structure; in other words Pr1-O bonds are, in average, under compressive stress and Pr2-O bonds are under tensile stress, giving rise to a structure with a slight instable character. The variability of oxidation states observed for Pr, able to adopt trivalent and tetravalent states with easiness can account for the clear deviation of Pr2 to an overbonded state. On the other hand, the valence of Mo ions is 5.5, noticeably below the 6+ oxidation state commonly found in many molybdates (for instance, CaMoO<sub>4</sub> with scheelite structure). Probably, the size of the tetrahedral coordination environment for Mo, conformed by the surrounding PrO<sub>8</sub> polyhedra in the fluorite substructure, is too large to allow Mo to satisfy its bonding requirement and to reach the hexavalent state. The fact that MoO<sub>4</sub> tetrahedra are isolated in this structure, having no common oxygen, prevents the electronic delocalization across the structure, since otherwise the intermediate valence state of Mo would lead to a black-colored compound with electronic conductivity. Transport measurements in sintered pellets show that Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub> is insulating at room temperature, with resistivities higher than  $10^8 \Omega$  cm. The variation of the electrical resistance with temperature was not measured.

The charge distribution observed in  $Pr_5Mo_3O_{16}$  is different from that obtained in other members of this structural family,  $CdTm_4Mo_3O_{16}$  [19] and  $CdY_4Mo_3O_{16}$  [23]. In these Cd-containing compounds, Cd and R=Y, Tm are occupying randomly the 8c sites (average oxidation state of 2.5+), whereas the 12e positions are uniquely occupied by  $R^{3+}$ , and Mo is, nominally, hexavalent. In the present case, Mo clearly exhibits a reduced oxidation state close to  $Mo^{5.5+}$ , whereas the bond valence obtained for Pr1 indicates a mixed  $Pr^{3+}-Pr^{4+}$  state, perhaps suggesting a partial charge disproportionation effect between Pr and Mo. An assumption of a purely trivalent Pr in  $Pr_5Mo_3O_{16}$  would give a nominal valence of 5.67+ for molybdenum.

The presence of structural stresses in the crystal structure can be quantified by means of the "Global Instability Index" [24], GII in Table 3, which is calculated from the differences between observed and expected valences for all the atoms in the asymmetric unit. It is a measure of the extent to which the BVS rule is violated over the whole structure. As suggested by Brown [24] and Armbruster et al. [25], GII values higher than 0.2 v.u. indicate the presence of intrinsic strains large enough to cause instability at room temperature. In the present case GII is 0.28 indicating a noticeable structural instability. This fact perhaps accounts for the stabilization of different crystal structures for other rare earths such as the  $\alpha$ - $R_2$ MoO<sub>6</sub> polytype, stable for R=Tb,..., Yb [14,16], or the  $\gamma$  polytype for La<sub>2</sub>MoO<sub>6</sub>.

Finally, it is worth extending the present results to the compound of nominal composition Ce<sub>2</sub>MoO<sub>6</sub>. As commented above, the similarity of the XRD diagrams of both oxides, including the superlattice reflections, suggests that both materials are isostructural, implying that the true composition of the Ce oxide may indeed be Ce<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub>. A study by X-ray spectroscopy of this material [4] already gave some insight on the coordination of the metal ions, in particular a distorted Ce-O<sub>8</sub> coordination environment and the tetrahedral location of the Mo ions. However, the description of Ce<sub>2</sub>MoO<sub>6</sub> with trivalent cerium and hexavalent Mo was inconsistent with previously proposed charge transfer models [15,26,27]. Moreover, as recognized by the authors of Ref. [4], the X-ray spectroscopy results lead to a quandary about the structure of Ce<sub>2</sub>MoO<sub>6</sub>, difficult to rationalize with a Ce<sup>3+</sup>/Mo<sup>6+</sup> cation distribution. The present structural model for Pr<sub>5</sub>Mo<sub>3</sub>O<sub>16</sub>, extrapolated to the Ce isomorph, could justify a Ce mixed valence, the presence of Mo in a reduced oxidation state and, plausibly, the occurrence of an internal charge transfer implying the reduction of Mo<sup>6+</sup> by Ce<sup>3+</sup> accounting for the unusual electronic properties of this material.

### 4. Conclusions

The  $\beta$ - $R_2$ MoO $_6$  polymorph has been investigated for R=Pr. The crystal structure, refined from NPD data, indeed corresponds to a  $Pr_5$ Mo $_3$ O $_{16}$  stoichiometry. It is a cubic superstructure of fluorite, formed by the long-range ordering of  $PrO_8$  and  $MoO_4$  coordination polyhedra. The two kinds of  $PrO_8$  scalenohedra share common edges and are interconnected by  $MoO_4$  tetrahedra by sharing corners with them. A bond-valence study indicates that the charge distribution observed in  $Pr_5$ Mo $_3$ O $_{16}$  is different from that obtained in other members of this structural family, such as  $CdTm_4$ Mo $_3$ O $_{16}$ , containing hexavalent molybdenum. In the present case, Mo ions clearly exhibit a reduced oxidation state close to  $Mo^{5.5+}$ , whereas the bond valence obtained for Pr1 indicates a mixed  $Pr^{3+}-Pr^{4+}$  state, perhaps suggesting a partial charge disproportionation effect between Pr and Pr0. Given the similarity of the XRD pattern of  $Pr_5$ Mo $_3$ O $_{16}$  with that published for  $Pr_2$ MoO $_6$ , it can be assumed that this latter oxide also

exhibits an analogous crystal structure and charge distribution, which needs to be further proved experimentally.

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